

1. If we consider QCD with *four* massless flavors, how many Goldstone bosons would arise in the breaking of chiral symmetry?
2. In the lecture, we encountered the relation (no summation over a)

$$\langle 0 | [Q_A^a, P^a(y)] | 0 \rangle = -\frac{1}{n_f} \langle 0 | \bar{\psi}(0) \psi(0) | 0 \rangle = -\langle 0 | \bar{u}(0) u(0) | 0 \rangle,$$

where n_f is the number of massless quarks, Q_A^a the axial charge, and

$$P^a(y) = \bar{\psi}(y) t^a \gamma_5 \psi(y).$$

Together with the Goldstone theorem, this implies that QCD has pseudoscalar Goldstone bosons, if the quark condensate is nonvanishing. Derive this relation using the canonical equal-time anti-commutation relations for the quark fields. The relation

$$[ab, cd] = a\{b, c\}d - ac\{b, d\} + \{a, c\}db - c\{a, d\}b$$

is useful.

3. The effective Lagrangian for the Goldstone bosons associated with spontaneous chiral symmetry breaking takes the form

$$\mathcal{L} = \frac{F^2}{4} \langle \partial_\mu U^\dagger \partial^\mu U \rangle + \mathcal{O}(p^4), \quad (1)$$

where $\langle \dots \rangle$ denotes the trace and the field $U(x)$ can be parameterized in terms of the Goldstone field π^a as

$$U(x) = \exp\left(\frac{2i}{F} \pi^a t^a\right),$$

The generators t^a correspond to the broken symmetries. Expand the Lagrangian to fourth order in the Goldstone fields π^a . For simplicity, work with two massless flavors, where $t^a = \sigma^a/2$.

4. Derive the Feynman rules for the Goldstone propagator and for the interaction of four Goldstone bosons.

5. Derive the Goldstone-boson scattering cross section, i.e. the cross section for the process

$$\pi^a(p_1) \pi^b(p_2) \rightarrow \pi^b(p_3) \pi^c(p_4).$$

6. To construct the effective Lagrangian, we added external sources for left-handed and right-handed vector currents to the QCD Lagrangian and transformed these as

$$\begin{aligned} r_\mu &\rightarrow V_R r_\mu V_R^\dagger - i(\partial_\mu V_R) V_R^\dagger \\ l_\mu &\rightarrow V_L l_\mu V_L^\dagger - i(\partial_\mu V_L) V_L^\dagger \end{aligned}$$

- (a) Show that with these transformations of the external fields the QCD Lagrangian becomes invariant under local chiral transformations $V_L(x)$ and $V_R(x)$.
- (b) Show that the leading-power chiral perturbation theory Lagrangian (1) is invariant under local transformations if we replace the regular derivatives by covariant one

$$i\partial_\mu \rightarrow iD_\mu = i\partial_\mu U + r_\mu U - U l_\mu.$$