

Color factors in Feynman diagrams

To evaluate Feynman diagrams in gauge theories, we need to simplify the color structure in Feynman diagrams. For color-singlet objects we encounter the objects

$$\delta^{ab}, f^{abc}, \text{tr} [t_R^{a_1} \cdot t_R^{a_2} \cdot \dots \cdot t_R^{a_n}]$$



in which all indices a, b, c, a_i are connected.

One can express these structures in terms of a few basic invariants of the underlying Lie group, namely the Casimir invariants

$$C_{R,R}(N) = \sum_{a_1, \dots, a_n} t_R^{a_1} \cdot t_R^{a_2} \cdot \dots \cdot t_R^{a_n} \text{tr} [t_R^{a_1} \cdot \dots \cdot t_R^{a_n}] \quad (*)$$

One can show that $C_{R, \mathbb{R}}(M)$ commutes with all generators t_R^a . If R is irreducible, this implies that $C_{R, \mathbb{R}}(M)$ is constant on R .

The most important invariant is the quadratic Casimir

$$C_R = C_{R, \mathbb{F}}(2) \frac{1}{T_{\mathbb{F}}} = t_R^a t_R^b \operatorname{tr}(t^a t^b) \cdot \frac{1}{T_{\mathbb{F}}} \\ = t_R^b t_R^b.$$

Let us verify that it commutes with all generators

$$\begin{aligned} [t_R^b \cdot t_R^b, t^a] &= t_R^b t_R^b t^a - t^a t_R^b t_R^b \\ &\quad + t_R^b t_R^a t_R^b - t_R^b t^a t_R^b \\ &= t_R^b i f^{abc} t_R^c - i f^{abc} t_R^c t_R^b \\ &= i f^{abc} (-t_R^b t_R^c - t_R^c t_R^b) = 0 \end{aligned}$$

$$\Rightarrow C_R \equiv C_R \cdot \mathbb{1}.$$

In the exercise class, we have evaluated C_F and C_A for $SU(N)$ and found

$$C_F = \frac{N^2 - 1}{2N} \quad ; \quad C_A = N$$

It turns out that up to two loops, the quadratic invariants are sufficient to evaluate any diagram. Beyond this higher invariants $C_{F,R}(N)$ are needed, but not all of these are independent.

First of all, it is sufficient to only consider symmetric trees. Antisymmetric pieces can be reduced using the group identity

$$[t_R^a, t_R^b] = i f^{abc} t_R^c.$$

One can also show that of the remaining invariants only r are independent. A detailed, algorithmic

way of computing color factors is presented in hep-ph/9802376.

A simple, but less general way of computing color structures in $SU(N)$ is to use the color Fierz identity

$$t_{ij}^a t_{ke}^a = \frac{1}{2} (d_{ie} d_{jk} - \frac{1}{N} d_{ij} d_{ke}) \quad (**)$$

and $f^{abc} = 2 \text{tr}([t^a, t^b] t^c)$. After eliminating the f^{abc} and using (**), all computations boil down to contractions of Kronecker δ 's.

This can easily be implemented into a computer code.